

Session of Monday 6 October at 16.00-17.30 in aud A315.

1. (a) Prove the following identity for a fermion propagator:

$$\frac{\not{p} + m}{p^2 - m^2 + i\varepsilon} = \frac{1}{2E_p} \sum_s \left[\frac{u(\mathbf{p}, s)\bar{u}(\mathbf{p}, s)}{p^0 - E_p + i\varepsilon} + \frac{v(-\mathbf{p}, s)\bar{v}(-\mathbf{p}, s)}{p^0 + E_p - i\varepsilon} \right]$$

where $E_p = +\sqrt{\mathbf{p}^2 + m^2}$. Note that the spinors are independent of p^0 .

Do a Fourier transform of the propagator from p^0 to time t , and give an interpretation of the two terms on the *rhs*.

- (b) Find the corresponding identity in terms of the LF spinors defined in eqs. (14) - (15) of the ‘‘LF Spinors and Polarization Vectors’’ file on the home page, where the ‘-’ momentum in the spinors is defined to be $p_{LF}^- \equiv (\mathbf{p}_\perp^2 + m^2)/p^+$.

Do a Fourier transform from p^- to x^+ and give your interpretation.

2. (a) Calculate the amplitude for a massless electron of momentum $p = (p^+, 0^-, \mathbf{0}_\perp)$ and positive helicity to radiate a photon with momentum $k = (xp^+, k_\perp^2/xp^+, \mathbf{k}_\perp)$ and helicity $\lambda = \pm 1$. You may use the LF spinors and polarization vectors given on the home page.
- (b) Discuss the dependence of the photon polarization on x . Find the relation between the squared amplitude and the splitting function $P_{\gamma/e}(x) = [1 + (1 - x)^2]/x$.
3. Protons may emit soft photons coherently, *i.e.*, photons with a wavelength longer than the proton radius, which couple to the total charge of the proton.
- (a) Estimate the range in transverse momentum and energy of photons coherently emitted by the 7 TeV protons in the LHC. Are they interesting for the LHC experiments?
- (b) Evaluate the off-shellness $p^2 - m^2$ of the emitting proton in the process $N(p) \rightarrow N(p') + \gamma(k)$. For a given $x = k^+/p^+$, estimate the range of \mathbf{k}_\perp in which the off-shellness is independent of the transverse momentum of the photon. From this, determine the effective radius of the photon cloud as a function of x . What is this radius (in fm) at the LHC for photons of 100 GeV?

$$1 \text{ a) } p^2 - m^2 + i\epsilon = (p^0 - E_p + i\epsilon)(p^0 + E_p - i\epsilon)$$

Check the common numerator on the rhs:

$$\frac{1}{2E_p} \left[(p^0 + E_p)(E_p \gamma^0 - \vec{p} \cdot \vec{\gamma} + m) + (p^0 - E_p)(E_p \gamma^0 + \vec{p} \cdot \vec{\gamma} - m) \right]$$

$$= p^0 \gamma^0 - \vec{p} \cdot \vec{\gamma} + m = \text{lhs.}$$

$$S(t, \vec{p}) = \int \frac{d^4 p}{(2\pi)^4} S(p, \vec{p}) e^{-itp^0} =$$

$$= \left[\theta(t) \int \frac{d^3 p}{(2\pi)^3} + \theta(-t) \int \frac{d^3 p}{(2\pi)^3} \right] S(p, \vec{p}) e^{-itp^0}$$

$$= -\frac{i}{2E_p} \sum_s \left[\theta(t) u(\vec{p}, s) \bar{u}(\vec{p}, s) e^{-itE_p} - \theta(-t) v(-\vec{p}, s) \bar{v}(-\vec{p}, s) e^{itE_p} \right]$$

⇒ First term is fermion propagating forward in t
 Second — " — anti — " — backward — " —

$$b) \text{ LF spinors: } \not{p}^- \equiv \frac{p_\perp^2 + m^2}{p^+}$$

$$U_{s,s} = \frac{1}{4} \not{x} (1 + s \gamma_5)$$

$$\sum_s u(p, s) \bar{u}(p, s) = \frac{1}{p^+} \sum_s (\not{p} + m) \underbrace{\tilde{\chi}(s) \chi^\dagger(s)}_{U_{s,s}} (\not{p} + m)$$

$$= \frac{1}{2p^+} (\not{p} + m) \not{x} (\not{p} + m)$$

$$= \frac{1}{2p^+} (\not{p} + m) \left[(-\not{p} + m) + 2 \underbrace{n \cdot p}_{p^+} \right] = \not{p} + m$$

16 cont.

$$\sum_s v(p, s) \bar{v}(p, s) = \frac{1}{2p^+} (\not{p} - m) \not{x} (\not{p} - m) = \not{p} - m$$

\Rightarrow As for ordinary spinors, but $p^- = \frac{p_\perp^2 + m^2}{p^+} \equiv p_{LF}^-$

$$\begin{aligned} \sum_s u(p, s) \bar{u}(p, s) &= \not{p}_{LF} + m = \frac{1}{2} p^+ \gamma^- + \frac{1}{2} p_{LF}^- \gamma^+ - \vec{p}_\perp \cdot \vec{\gamma}_\perp + m \\ &= \frac{1}{2} p^+ \gamma^- + \frac{1}{2} p^- \gamma^+ - \vec{p}_\perp \cdot \vec{\gamma}_\perp + m + \frac{1}{2} (p^- - p_{LF}^-) \gamma^+ \end{aligned}$$

$$p^2 - m^2 + i\varepsilon = p^+ p^- - p_\perp^2 - m^2 + i\varepsilon = p^+ (p^- - p_{LF}^-) + i\varepsilon$$

$$\frac{\not{p} + m}{p^2 - m^2 + i\varepsilon} = \sum_s \frac{u(p, s) \bar{u}(p, s)}{p^+ (p^- - p_{LF}^-) + i\varepsilon} + \frac{1}{2p^+} \gamma^+$$

$$\begin{aligned} \int \frac{dP^-}{2\pi} S(p^+, p; \vec{p}_\perp) e^{-\frac{i}{2} x^+ P^-} &= -i \Theta(x^+) \Theta(p^+) \frac{1}{p^+} \sum_s u(p, s) \bar{u}(p, s) e^{-ix^+ p_{LF}^- / 2} \\ &+ i \Theta(-x^+) \Theta(-p^+) \frac{1}{p^+} \sum_s u(p, s) \bar{u}(p, s) e^{-ix^+ p_{LF}^- / 2} \\ &+ \delta(x^+ / 2) \gamma^+ / 2p^+ \end{aligned}$$

$$= -\frac{i}{|p^+|} \Theta(x^+ p^+) \sum_s u(p, s) \bar{u}(p, s) e^{-ix^+ p_{LF}^- / 2} + \frac{\gamma^+}{p^+} \delta(x^+)$$

* Fermions with $p^+ > 0$ propagate forward in x^+
 $p^+ < 0$ propagate backward in x^+

* There is an "instantaneous" (in x^+) component, which may be thought of as the backward-propagating (in t) anti-fermion boosted to the ∞ momentum frame.

2a) Helicity is conserved
for $m=0$ fermions



$$iA_\lambda = \bar{u}(p', +) (-ie \not{\epsilon}_\lambda^*(k)) u(p, +)$$

$$p = (p^+, 0, \vec{0}_\perp) = \frac{1}{2} p^+ \tilde{n}$$

$$u(p, +) = \frac{1}{\sqrt{p^+}} \not{\epsilon} \tilde{\chi}(+)$$

Note: $p'^- = \frac{P_\perp^2}{p'^+}$ in LF spinor

$$\neq p^- - k^-$$

$$A_\lambda = -\frac{e}{\sqrt{p^+ p'^+}} \text{Tr} \left[\not{p}' \not{\epsilon}_\lambda^*(k) \not{\epsilon} \underbrace{\tilde{\chi}(+) \chi^+(+)}_{U_{++}} \right]$$

$$U_{++} = \frac{1}{4} \not{n} (1 + \gamma_5) \quad \left(\begin{array}{l} n \cdot p = p^+ \\ n \cdot n = 0 \end{array} \right)$$

$$= -\frac{e}{4\sqrt{p^+ p'^+}} \text{Tr} \left[\not{p}' \not{\epsilon}_\lambda^* \not{\epsilon} \not{n} (1 + \gamma_5) \right] \quad (\lambda = \pm 1)$$

$$\epsilon_\lambda = e_\lambda - \frac{e_\lambda \cdot k}{k^+} n \quad ; \quad e_\lambda = (0, 0, \vec{e}_\lambda) \quad ; \quad e_\lambda = -\frac{1}{\sqrt{2}} (\lambda, i)$$

$$(\lambda = \pm 1)$$

$$e_\lambda \cdot k = \lambda \frac{k_\perp}{\sqrt{2}} e^{i\lambda\varphi} \quad ; \quad \vec{k}_\perp = k_\perp (\cos\varphi, \sin\varphi)$$

$$n \cdot \epsilon_\lambda^* = 0$$

$$p \cdot \epsilon_\lambda^* = -\frac{\lambda}{\sqrt{2}} k_\perp e^{-i\lambda\varphi}$$

$$p' \cdot \epsilon_\lambda^* = -k \cdot \vec{e}_\lambda^* - \frac{e_\lambda^* \cdot k}{x p^+} (1-x) p^+ = -\frac{e_\lambda^* \cdot k}{x} = p \cdot \epsilon_\lambda^*$$

$$A_\lambda = -\frac{e}{p^+ \sqrt{1-x}} \left\{ \left[p^+ (p' \cdot \epsilon_\lambda^*) + p'^+ (p \cdot \epsilon_\lambda^*) \right] + \frac{p^+}{8} \text{Tr} \left[\not{p}' \not{\epsilon}_\lambda^* \not{n} \not{n} \gamma_5 \right] \right\}$$

Use identity: $\not{n} \gamma_5 e_\lambda = \lambda \not{n} e_\lambda \quad (8)$

$$e_\lambda^* = -e_{-\lambda}$$

$p' \rightarrow -k$ in trace (only \perp contrib.)

2a cont.)

$$\begin{aligned} \text{Tr} &= +\text{Tr} [k \tilde{\gamma} \not{x} \gamma_5 \not{e}_{-\lambda}] = +\lambda \text{Tr} [k \tilde{\gamma} \not{x} \not{e}_{-\lambda}] \quad |n \cdot \tilde{n} = 2 \\ &= +8\lambda k \cdot e_{-\lambda} = -\frac{8}{\sqrt{2}} k_{\perp} e^{-i\lambda\varphi} \quad (\lambda^2 = 1) \end{aligned}$$

$$A_{\lambda} = + \frac{e}{\sqrt{1-x}} \frac{k_{\perp} e^{-i\lambda\varphi}}{\sqrt{2} x} \left[\lambda + (1-x)\lambda + x = (1+\lambda) - (1-\lambda)(1-x) \right]$$

b)

$$\frac{A_{-1}}{A_{+1}} = -(1-x) \quad \begin{cases} \rightarrow 0 \text{ for } x \rightarrow 1: \gamma \text{ carries } e \text{ spin direction} \\ \rightarrow -1 \text{ for } x \rightarrow 0: \gamma \text{ unpolarized} \end{cases}$$

$$\sum_x |A_{\lambda}|^2 = \frac{2e^2 k_{\perp}^2}{x(1-x)} \frac{1 + (1-x)^2}{x} = \frac{2e^2 k_{\perp}^2}{x(1-x)} P_{e \rightarrow \gamma}(x)$$

3. $P \rightarrow p + \gamma$ coherently

$$\Rightarrow P_\gamma \lesssim \frac{1}{r_p} \approx \frac{1}{1 \text{ fm}} \approx 200 \text{ MeV}$$

a) At the LHC: $P_\perp \lesssim 200 \text{ MeV}$

$$P_{\parallel} \lesssim 0.2 \text{ GeV} \cdot \frac{E_p}{m_p} = 0.2 \cdot \frac{7 \cdot 10^3}{.94} = 1.5 \text{ GeV}$$

$$x = \frac{E_\gamma}{E_p} \approx \frac{1}{r_p} \frac{E_p}{m_p} \frac{1}{E_p} = \frac{1}{m_p r_p} \approx 0.2$$

Yes: 1.5 GeV γ + 7 TeV p gives

$$E_{\text{cm}} \approx \sqrt{4 \cdot 1.5 \cdot 7} = 6.5 \text{ TeV} \text{ is interesting!}$$

$$b) (p' + k)^2 - m_p^2 = 2 p' \cdot k = p'^+ k^- + p'^- k^+ + 2 k_\perp^2$$

$$= (1-x) p^+ \frac{k_\perp^2}{x p^+} + \frac{k_\perp^2 + m_p^2}{(1-x) p^+} x p^+ + 2 k_\perp^2 =$$

$$= \frac{x}{1-x} m_p^2 + k_\perp^2 \left[\frac{1-x}{x} + \frac{x}{1-x} + 2 \right] = \frac{(1-x)^2 + x^2 + 2x(1-x)}{x(1-x)} = \frac{1}{x(1-x)}$$

$$= \frac{1}{x(1-x)} \left[k_\perp^2 + (x m_p)^2 \right] \text{ which is insensitive to } k_\perp \text{ for}$$

$$k_\perp \lesssim x m_p \Rightarrow r_\gamma \approx \frac{1}{k_\perp} = \frac{1}{x m_p}$$

$$x = \frac{0.1}{7} \Rightarrow r_\gamma = \frac{7}{0.1 \cdot 0.94 \text{ GeV}} \cdot 0.2 \text{ fm GeV} = 15 \text{ fm}$$

Note: $r_\gamma \rightarrow \infty$ for $x \rightarrow 0$: EM has ∞ range